

Article

Fluid–Structure Interaction Modeling Applied to Peristaltic Pump Flow Simulations

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Received: 29 March 2019; Accepted: 24 June 2019; Published: 9 July 2019

Abstract: In this study, fluid–structure interaction (FSI) modeling was applied for predicting the fluid flow in a specific peristaltic pump, composed of one metallic roller and a hyperelastic tube pumping a viscous Newtonian fluid. Hyperelastic material dynamics and turbulence flow dynamics were coupled in order to describe all the physics of the pump. The commercial finite element software ABAQUS 6.14 was used to investigate the performance of the pump with a 3D transient model. By using this model, it was possible to predict the von Mises stresses in the tube and flow fluctuations. The peristaltic pump generated high pressure and flow pulses due to the interaction between the roller and the tube. The squeezing and relaxing of the tube during the operative phase allowed the liquid to have a pulsatile behavior. Numerical simulation data results were compared with one cycle pressure measurement obtained from pump test loop data, and the maximum difference between real and simulated data was less than 5%. The applicability of FSI modeling for geometric optimization of pump housing was also discussed in order to prevent roller and hose parts pressure peaks. The model allowed to investigate the effect of pump design variations such as tube occlusion, tube diameter, and roller speed on the flow rate, flow fluctuations, and stress state in the tube.

Keywords: peristaltic pump; hyperplastic models; turbulence models; pulsatile flow

1. Introduction

The peristaltic pump is a pumping device that is based on the peristalsis principle. The specific pump of this work was composed of a cylindrical rotor that squeezed and relaxed an elastic pipe where the fluid flows. In addition, during the operative phase, the pipe was submitted to high stress values, which might have caused such high deformation on the pipe that it affected the liquid's transport. Experimental analysis could be used for developing pressure peak, pump optimizing, and for selecting optimal piping material. A feasible alternative to experimental analysis is the numerical simulation where important design variables could be systemically studied. From experimental sensitivity analysis, the Poisson's coefficient of the tube material was very important. The Poisson coefficient (also called transversal contraction coefficient) is a property of each material, depending on the temperature, which measures the degree to which the material sample shrinks or dilates transversely in the presence of a longitudinal unidirectional stress.

The values of the Poisson's coefficient for materials in nature are between 0 and 0.5; the value 0.5 corresponds to a virtually incompressible material (the rubber, for example, has values close to 0.5) which therefore has an infinite volumetric expansion module.

The peristaltic pump tubes are usually made of materials such as silicone, PVC, and other polymers that can offer, in addition to the aforementioned mechanical characteristics, also a remarkable chemical compatibility with solvents, acids, inks, paints, etc. Elastomers reinforced with textile inserts are used to guarantee the toughness of the pipe even at high pressures and for a long time.

The peristaltic pump is used in all processes in which the treated fluid has to get far from the pump components for safety reasons against contamination, such as in the food, pharmaceutical, and medical industries, or when products with solids in suspension must be dosed without being damaged. This type of pump allows to accurately calibrate the flow rate and keep it constant. These characteristics make it ideal where precision on the dosages and repeatability of the measurement are indispensable requirements such as in chemical laboratories, in bio-research, etc.

The need for a more accurate control of fluid flow for applications in Medicine and Engineering has increased the interest in performing numerical simulations. The control of the fluid flow derives from the fluid-induced deformation or the boundary moving induced deformation. The simple hose of the peristaltic pump's structure can perform control and non-blocking high pressure, or even shock pressure, pumping throughout pipe deformations caused by the pump actuator parts. Furthermore, the high deformation may cause permanent damage for the hyperelastic pipe that affects liquid transport—optimal tubing material can be developed by pressure peak of the pumps.

The design of a peristaltic pump requires an accurate modelling [\[1](#page-9-0)[,2\]](#page-9-1). The pump could be considered as a multibody model [\[3](#page-9-2)[–6\]](#page-9-3) and used in engineering applications [\[7](#page-9-4)[–9\]](#page-9-5) for dynamic simulations [\[10,](#page-9-6)[11\]](#page-9-7). Active vibration control is another aspect to consider during the design phase [\[12\]](#page-9-8). Dynamic analysis [\[13\]](#page-9-9) of each considered problem has some peculiar difficulties. In [\[14\]](#page-9-10), the authors describe the influence of peristaltic fluid flow (blood) on the elastic boundaries (the blood vessel); in this case FSI problem is approached [\[15–](#page-9-11)[17\]](#page-9-12).

To describe an accurate model behavior of peristaltic pumps, fluid and solid models have to be merged—fluid flow is caused by pipe squeezing. The deformation of the tube is strongly coupled to the fluid flow, so an FSI analysis is required; then FSI modeling is mainly spread for pulsatile flow induced by moving domain. The data results are compared to the relative fixed boundary domain flow problems because of the small deformation of the domain, contrary to the hose system. Elabbasi [\[18\]](#page-9-13) investigated the performance of a 180-degree rotary peristaltic pump with two metallic rollers by using 3D FSI with COMSOL Multiphysics. The model merges the fluid and structural parts in a monolithic way—the structure and the fluid parts are respectively described by Mooney-Rivlin hyperelastic material model and laminar flow models. However, the authors do not give any details on mathematics and they exclude any turbulence phenomena for the fluid. Zhou [\[19\]](#page-9-14) described a 2D CFD tool for studying peristaltic pumps using ANSYS Fluent. This program code allowed to approach the Computational Mechanics [\[20–](#page-10-0)[22\]](#page-10-1), and it was also used to perform the control of nonlinear vibrations [\[23](#page-10-2)[–25\]](#page-10-3) for the machine design [\[26](#page-10-4)[–28\]](#page-10-5). Further, it allowed the determination of the algorithms for computing the modal parameters of linear mechanical systems [\[29](#page-10-6)[,30\]](#page-10-7).

Therefore, this study aimed to develop a mathematical and numerical 3D transient FSI model for an industrial peristaltic pump. The main design variables governing pump performance were the pump speed, inner tube diameter, tube material, degree of tube occlusion, and back pressure. The model has been implemented in commercial software ABAQUS and catches the pipe deformation by a Neo-Hookean hyperelastic model and the fluid flow by the RNG κ - ϵ turbulence model. The mesh equation was described as a dynamic harmonic meshing function [\[31\]](#page-10-8). Therefore, through this study, it was seen that it is possible to simulate the operating conditions of a peristaltic pump and to obtain a numerical model of the peristaltic pump considered that allowed to perform a better study of the working phase aimed to improve and to optimize the operating characteristics of the peristaltic pump considered.

2. Geometry and Model 2. Geometry and Model

2.1. Model Geometry and Finite Element Mesh 2.1. Model Geometry and Finite Element Mesh

In the peristaltic pump, a rotating roller squeezes a resilient pipe and the fluid is pushed forward In the periodicity pump, a rotating roller squeezes a resilient pipe and the fluid is pushed forward.
It through the hose pipe with compression and relaxation phases, see Figure [1.](#page-2-0) The rotation period of the through the hose pipe with compression and relaxation phases, see Figure 1. The rotation period of the roller and distance of roller axis from pipe axis affects the fluid flow along the hose pipe. the roller and distance of roller axis from pipe axis affects the fluid flow along the hose pipe.

Figure 1. Schematic illustration of peristaltic pump. **Figure 1.** Schematic illustration of peristaltic pump.

The FSI model of the peristaltic pump consisted of a cylindrical rigid roller, a round helix hose The FSI model of the peristaltic pump consisted of a cylindrical rigid roller, a round helix hose pipe, and a cylindrical rigid frame that covered the device and bounded the pipe movement during pipe, and a cylindrical rigid frame that covered the device and bounded the pipe movement during the squeezing phase. The features of the pipe made essential a 3D analysis of the phenomenon. the squeezing phase. The features of the pipe made essential a 3D analysis of the phenomenon.

Radius of the rigid roller was 132.5 mm, and the pipe had a thickness of 15 mm, an internal Radius of the rigid roller was 132.5 mm, and the pipe had a thickness of 15 mm, an internal radius of 32 mm, and an external one of 146 mm. The inertial property of the roller was neglected.

The structural part was meshed using 145 rectangular linear shell elements for the fixed frame, The structural part was meshed using 145 rectangular linear shell elements for the fixed frame, 40058 tetrahedral quadratic elements for the hose pipe Ω^{*s*}, 208 hexahedral linear elements for the roller Ω^r , se[e F](#page-2-1)igure 2. The reference fluid domain, Ω^f_0 , was meshed using 46994 tetrahedral elements, as in Figure 3. The boundary layer had size 9% of the global mesh size and it consisted of 5 layers. \int_0^7 , was meshed using 46994 tetrahedral elements, as in

Figure 2. Mesh of the structural part.
 Figure 2. Mesh of the structural part.

Figure 3. Mesh of the structural part. **Figure 3.** Mesh of the structural part.

2.2. Mathematical Model The height of the wall boundary layer during the squeezing was observed as having a significant but the element nodes on the common boundary coincided. t ransient analysis. A turbulence model is needed to imitate actual flows in hose pumps and in order α effect on the velocity field convergence. The fluid and structure geometries were separately meshed

2.2. Mathematical Model **is a computation of computation** α computer between the computational model is a computational model in the computational model is a computational model in the computational model in the computa

cost and the efficient of the results for FSI phenomena in FSI phenom The finite deformations of the pipe in peristalsis phenomena rise up turbulent structures in a
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model the fluid sub-problem. A RANS-based model is a compromise between the computational cost and the efficient of the results for FSI phenomena [32-[36\]](#page-10-10). transient analysis. A turbulence model is needed to imitate actual flows in hose pumps and in order to

థ The fluid is described in an Arbitrary Lagrangian Eulerian framework. The ALE derivative The fluid is described in an Arbitrary Lagrangian Eulerian framework. The ALE derivative
induced by a deformation of fluid domain, $\phi : \Omega_0^f \times I \to R^3$ for a scalar function $f : \Omega_0^f \times I \to R$ is Formed as $D^{\phi}f = \partial f - \partial \phi$ \int_0^f \times *I* \rightarrow *R*³ for a scalar function $f : \Omega_0^f$ $\frac{f}{0} \times I \to R$ is defined as

$$
\frac{D^{\phi}f}{Dt} := \frac{\partial f}{\partial t} + \frac{\partial \phi}{\partial t} \nabla f \tag{1}
$$

 $\ddot{}$ by the RNG κ – ϵ turbulence model in an ALE framework The deformed fluid domain is defined as $\forall t \in I$ as $\Omega_t^f := \phi(\Omega_0^f, t)$. The fluid m $t_t^f := \phi(\Omega_0^f)$ $\int_0^f f(x) dx$. The fluid motion was described

$$
\begin{cases}\n\rho^f \frac{D^{\phi} U^f}{Dt} + \rho^f (U^f - w) \nabla U^f = -\nabla P^f + (\mu^f + \mu_t) \Delta U^f - \frac{2}{3} \nabla k \\
\rho^f \frac{D^{\phi} k}{Dt} + \rho^f (U^f - w) \cdot \nabla k = (\mu^f + \frac{\mu_t}{\sigma_k}) \Delta k - \rho^f \epsilon + P_k \\
\rho^f \frac{D^{\phi} \epsilon}{Dt} + \rho^f (U^f - w) \cdot \nabla \epsilon = (\mu^f + \frac{\mu_t}{\sigma_{\epsilon}}) \Delta \epsilon - C_{2\epsilon} \rho^f \frac{\epsilon^2}{k} \epsilon + \frac{\epsilon}{k} C_{1\epsilon} P_k - R_{\epsilon}\n\end{cases} (2)
$$

where U^f is the mean velocity field, P^f is the mean pressure, k is the turbulent kinetic energy, ϵ is the turbulent dissipation energy, ρ^f is the fluid density, μ^f is the fluid dynamic viscosity, P_k is the rate of production of turbulence, R_{ϵ} is the rate of dissipation, and μ_t is the turbulent dynamic viscosity. The coefficient σ_k , σ_ϵ , $C_{1\epsilon}$, and $C_{2\epsilon}$ are closure coefficients. A linear elastic model is not able to catch the finite deformation of the pipe [\[37\]](#page-10-11).

All of that is also considered in the design of complex mechanical systems [\[38](#page-10-12)[–40\]](#page-10-13) that involved models with hysteresis [\[41](#page-10-14)[,42\]](#page-10-15) aimed to solve thermodynamic problems [\[43\]](#page-10-16) and to perform dynamic analysis [\[44–](#page-11-0)[46\]](#page-11-1) and optimization [\[47](#page-11-2)[–49\]](#page-11-3) evaluating also the structural displacements.

The pipe displacement was determined by a Neo-Hookean hyperelastic model in a Lagrangian framework

$$
\rho^s \frac{d^2 \eta^s}{dt^2} + \rho^s \alpha \frac{d\eta^s}{dt} - \nabla \cdot \widetilde{T}^s - \beta \frac{d}{dt} \nabla \cdot \widetilde{T}^s = 0 \tag{3}
$$

where η^s is the displacement field, ρ^s is the pipe density, \overline{T}^s is the stress tensor, α is the mass damping coefficient, and β is the stiffness damping coefficient.

The displacement field of the roller φ_r is computable by simple rigid body equations by fixing the rotation period *T*. The interaction between the pipe and the roller is frictionless—no dissipative factors are taken on count. Furthermore, the roller is not subjected by the counter force of the pipe. Furthermore, the problem involved an ALE smooth mesh problem for the resolution of geometry

ALE mesh problem in
$$
\Omega_t^f
$$
: $\begin{cases} \Delta \varnothing = 0 \\ w = \frac{\partial \phi}{\partial t} \end{cases}$
The boundary conditions are:

The boundary conditions are:

For the fluid part: No-flux condition on the ends;

For the solid part: No-stress condition on the free surface.

The coupling conditions between the fluid and structure on the movable interface were kinematic, i.e., continuity of velocity, and dynamic, i.e., continuity of nominal stress.

For the time discretization a Backward-Euler scheme was used.

The working fluid was water with density $\rho^f = 10^3$ kg/m³ and dynamic viscosity $\mu_t = 10^{-3}$ Pa s. For the fluid domain stress free conditions at inlet and outlet were imposed. The material parameters of the hyperelastic pipe are given in Table [1.](#page-4-0) The pipe-roller and pipe-frame contacts were assumed to be frictionless. Self-contact of the pipe was excluded.

[kg/m ³]	[MPa]		c_{10} [MPa]	$[MPa^{-1}]$	a $[s^{-1}]$	[s]
970	37.9	0.48	127	48	0.4	U. 1

Table 1. Material parameters of the pipe.

The ρ *^s* parameter is the density, the *E* parameter is the Young's modulus, the ν parameter is the Poisson ratio, the *C*¹⁰ parameter weights the deformation of the body, the *D*¹ parameter weights the effect of compressibility along the motion, the α parameter is the mass damping coefficient and β is the stiffness damping coefficient. The parameters are computed as in Bergstrom [\[50\]](#page-11-4).

The ν value considered was around 0.4; 0.5 corresponds to a virtually incompressible material.

2.3. Simulation Set-Up

In the simulations ABAQUS 6.14.1 software was used. The fluid part was solved using ABAQUS/CFD and the structural part using ABAQUS/Standard.

At the beginning of the simulation, the roller was not touching the pipe. The simulation was initiated by compressing the roller, using uniform linear motion against the pipe (see Figure [4\)](#page-5-0). The compression phase was completed after 5 s, see Figure [5.](#page-5-1) The compression reduced the internal diameter of the pipe about 87%, from 32 mm to 4 mm.

Figure 4. Phase 1 of simulation. **Figure 4.** Phase 1 of simulation.

Figure 5. Phase 2 of the simulation. **Figure 5.** Phase 2 of the simulation. **Figure 5.** Phase 2 of the simulation.

3. Results and Discussion 3. Results and Discussion

3.1. Numerical Simulation Tests 3.1. Numerical Simulation Tests

In Figures [6](#page-5-2) and [7,](#page-6-0) the output velocity diagram obtained during the numerical simulation are reported. The maximum fluid velocity at the end of the first phase of the simulation was around the velocity was pulsatile. 0.05 m/s. The max[im](#page-6-0)um velocity value was around 0.67 m/s, and in Figure 7, it is possible to note that

Figure 6. Phase 2 of the simulation; velocity diagram.

Figure 7. Output velocity obtained during the numerical simulation.

The compression significantly affected the time step. The more the pipe was squeezed, the more the time step of the time discretization scheme decreased. During the squeezing, the time step of the Backward Euler scheme decreased from 10−1 s to 10−3 s. Backward Euler scheme decreased from 10−¹ s to 10−³ s. Backward Euler scheme decreased from 10−1 s to 10−3 s. Backward Euler scheme decreased from 10−1 s to 10−3 s.

Furthermore, when the pipe was compressed at both ends, the time step decreased to 10^{-6} s. After the initial compression, the rolling process was activated by rotating the roller with prescribed velocity, as illustrated in Figure [8.](#page-6-1) The period of rotation was 1.1 s.

Figure 8. Rotation of the roller with period of around 1.1 s.

The von Mises stress of the pipe reaction in the roller stress of the roller squeezed the pipe. When $\frac{1}{2}$ the roller moved on, the Mises stress decreased with hyperbolic behavior, see Figur[e 9](#page-6-2). The von Mises stress of the pipe reached 1.4×10^7 Pa when the roller squeezed the pipe. When

Figure 9. Von Mises stress at tube outlet. **Figure 9.** Von Mises stress at tube outlet. **Figure 9.** Von Mises stress at tube outlet. **Figure 9.** Von Mises stress at tube outlet.

 F_{F} $\frac{1}{2}$ from the simulation, output pressure decreased with hyperbolic behavior after it had reached with hyperbolic behavior after it has reached with hyperbolic behavior after it had reached with $\frac{1}{2}$ a maximum values, Fi[gure](#page-7-0) 10. From the simulation, output pressure decreased with hyperbolic behavior after it had reached the linear related U

Figure 10. Pressure data at output of numerical simulation for the considered peristaltic pump. **Figure 10.** Pressure data at output of numerical simulation for the considered peristaltic pump. Figure 10. I fessure that at output of numerical simulation for the considered peristant pump.

For the numerical simulation, an Arbitrary-Lagrangian-Eulerian mesh transformation was used—the squeeze of the fluid mesh was set on the basis on the numerical consideration.

For the specific problem, robust algorithms are needed in order to numerically simulate the real For the specific problem, robust algorithms are needed in order to numerically simulate the real dosing phenomena where the squeezing holds a self-interaction between the internal hose pump dosing phenomena where the squeezing holds a self-interaction between the internal hose pump wall. *3.2. Experimental Tests* The future aims are to reduce the squeezed diameter close to zero.

3.2. Experimental Tests

In Figure [11,](#page-7-1) the experimental data detected for the case study peristaltic pump were reported with one cycle pressure measurement. The data reported in Figure [11](#page-7-1) were the mean values of data results of experimental tests repeated three times, and the difference between the maximum and minimum values obtained was less than 3%. When the roller squeezed both pipe ends, a pressure drop state for the average pressure was detected as in Figure [11.](#page-7-1) The roller positions were analyzed from video afterwards. Pressure and video data results were synchronized to the same time scale. Although the experimental data was collected at sparse time intervals (200 ms), the pump operation became clear when the peaks were formed.

Figure 11. Experimental pressure data at output of a peristaltic pump. **Figure 11.** Experimental pressure data at output of a peristaltic pump.

The maximum pressure peak occurred in a cyclic movement just to the position when both hoses were squeezed by roller. At other roller positions, such pressure differences are not observed. Further, in Figure [11](#page-7-1) we can note a periodic behavior at longer time scale.

4. Discussion

The numerical data results in Figure [10](#page-7-0) and real data in Figure [11](#page-7-1) are very similar along one revolution of about 1 s; the maximum difference evaluated was lower than 5%.

For the elastic pipes, during the simulations, it was possible to detect different pulsatile phases in the fluid along the motion and to describe the behavior of the movement of the wall.

Other elements of simulation were:

- 1. The absence of pressure counter in the model allowed a decrease in the pressure output. In the experimental case, there was a valve that generated pressure counter against the back flow.
- 2. Different pipe material: The material parameter, as underlined in Equation (2), influences the fluid flow.
- 3. The internal diameter when the roller squeezed the pipe was reduced by about 87%, (from 32 mm to 4 mm).

Other elements can be added in order to get a more realistic model:

- Dissipative friction effect between pipe and roller;
- Roller as deformable material.

From the data results, we found out that it was possible to simulate the singular position in the pipe. Further, when the pipe compressed the inlet and outlet, the pressure had discontinuity that was detected by the numerical model. This discontinuity is due to the geometrical shape of the pipe. As one major advantage, this model, although having minor deficiencies, enables to design new generation peristaltic pump geometries that can have minimal pressure differences during pump operation.

5. Conclusions

The aim of this study was the development of a powerful fluid simulation tool that allowed to optimize the peristaltic pump geometry. The numerical model was able to predict the pump behavior during the operative phase. The simulation data results were compared to the experimental ones. The simulated and experimental cyclic pressure data agreed well with each other; simulated pressure magnitudes were matched, and the maximum difference evaluated between numerical and experimental data obtained was lower than 5%. Therefore, another possible tool application was the geometric optimization of pump housing, roller, and hose parts for developing pressure peak optimized pumps.

Further, in this report, a study of advanced mathematical tools for defective boundary conditions problems for engineering approximations in fluid dynamics was performed. A numerical algorithm was built for the resolution of a laminar flow problem with mean pressure (defective) inlet condition and for the resolution of a steady RANS-based turbulent model with unknown Boundary-Layer height. Therefore, pulsatile flows were examined with FSI and defective analysis by using the commercial software ABAQUS, which is based on a Finite Element method, and then the simulation of the hose pump was performed. For FSI problems, the fluid flow was described by a RANS-based model called RNG $k - \epsilon$ in an ALE framework. The FSI-modeling was used for the pumping simulation of the main pump parts interaction, and very good results were obtained.

The FSI model simulations have not been reported vastly in scientific journals and conferences. The few articles do not contain detailed descriptions either of the models or of the results.

As one major goal of this study, this model, although having minor deficiencies, will enable the design of new generation peristaltic pump geometries that can have minimal pressure differences during pump operation.

Author Contributions: All the authors equally contributed to carry out the research whose results are reported in this work.

Funding: This research received no external funding.

Conflicts of Interest: The authors declare no conflict of interest.

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